

Getting the Darcy law by time rescaling: a quick way to derive models for porous media from "first principles"

Yann Brenier

CNRS, Centre de mathématiques LAURENT SCHWARTZ
Ecole Polytechnique, FR-91128 Palaiseau

Modeling and Simulation in Porous Media

INRIA, Rocquencourt, 8-9 Dec. 2014

Deriving the Darcy law from the Euler equations

In 1755/57, Euler introduced the first "field theory" in Physics, and, at the same time, the first nonlinear PDE ever written

$$\partial_t \rho + \operatorname{div}(\rho \mathbf{v}) = 0, \quad \partial_t(\rho \mathbf{v}) + \operatorname{div}(\rho \mathbf{v} \otimes \mathbf{v}) = -\operatorname{grad} p$$

where $(\rho, p, \mathbf{v}) \in \mathbb{R}^{1+1+3}$ are the density, pressure and velocity fields of a fluid and ρ is assumed to be a given function of ρ .



XXI. Nous n'avons donc qu'à équaler ces forces accélératrices avec les accélérations actuelles que nous venons de trouver, & nous obtiendrons les trois équations suivantes :

$$P - \frac{1}{q} \left(\frac{dp}{dx} \right) = \left(\frac{du}{dt} \right) + u \left(\frac{du}{dx} \right) + v \left(\frac{du}{dy} \right) + w \left(\frac{du}{dz} \right)$$

$$Q - \frac{1}{q} \left(\frac{dp}{dy} \right) = \left(\frac{dv}{dt} \right) + u \left(\frac{dv}{dx} \right) + v \left(\frac{dv}{dy} \right) + w \left(\frac{dv}{dz} \right)$$

$$R - \frac{1}{q} \left(\frac{dp}{dz} \right) = \left(\frac{dw}{dt} \right) + u \left(\frac{dw}{dx} \right) + v \left(\frac{dw}{dy} \right) + w \left(\frac{dw}{dz} \right)$$

Si nous ajoutons à ces trois équations premièrement celle, que nous a fournie la considération de la continuité du fluide :

$$\left(\frac{dq}{dt}\right) + \left(\frac{d.qu}{dx}\right) + \left(\frac{d.qv}{dy}\right) + \left(\frac{d.qw}{dz}\right) = 0.$$

Si le fluide n'étoit pas compressible, la densité q seroit la même en Z , & en Z' , & pour ce cas on auroit cette équation :

$$\left(\frac{du}{dx}\right) + \left(\frac{dv}{dy}\right) + \left(\frac{dw}{dz}\right) = 0.$$

qui est aussi celle sur laquelle j'ai établi mon Mémoire latin allégué ci-dessus.

An unusual but simple trick to get **the Darcy law** out of the Euler equations consists in rescaling time in a nonlinear way as follows

$$t \rightarrow \tau = t^2/2, \quad (\rho, v)(t, x) \rightarrow (\rho(\tau, x), \tau' v(\tau, x)), \quad \tau' = \frac{d\tau}{dt} = t$$

An unusual but simple trick to get **the Darcy law** out of the Euler equations consists in rescaling time in a nonlinear way as follows

$$t \rightarrow \tau = t^2/2, \quad (\rho, \mathbf{v})(t, \mathbf{x}) \rightarrow (\rho(\tau, \mathbf{x}), \tau' \mathbf{v}(\tau, \mathbf{x})), \quad \tau' = \frac{d\tau}{dt} = t$$

$$\partial_t \rho + \operatorname{div}(\rho \mathbf{v}) = \mathbf{0}, \quad \partial_t(\rho \mathbf{v}) + \operatorname{div}(\rho \mathbf{v} \otimes \mathbf{v}) = -\operatorname{grad} \mathbf{p} : \rightarrow$$

An unusual but simple trick to get **the Darcy law** out of the Euler equations consists in rescaling time in a nonlinear way as follows

$$t \rightarrow \tau = t^2/2, \quad (\rho, \mathbf{v})(t, \mathbf{x}) \rightarrow (\rho(\tau, \mathbf{x}), \tau' \mathbf{v}(\tau, \mathbf{x})), \quad \tau' = \frac{d\tau}{dt} = t$$

$$\partial_t \rho + \operatorname{div}(\rho \mathbf{v}) = \mathbf{0}, \quad \partial_t(\rho \mathbf{v}) + \operatorname{div}(\rho \mathbf{v} \otimes \mathbf{v}) = -\operatorname{grad} \mathbf{p} : \rightarrow$$

$$(\partial_\tau \rho + \operatorname{div}(\rho \mathbf{v})) \tau' = \mathbf{0}, \quad \tau'' \rho \mathbf{v} + (\tau')^2 [\partial_\tau(\rho \mathbf{v}) + \operatorname{div}(\rho \mathbf{v} \otimes \mathbf{v})] = -\operatorname{grad} \mathbf{p}$$

An unusual but simple trick to get **the Darcy law** out of the Euler equations consists in rescaling time in a nonlinear way as follows

$$t \rightarrow \tau = t^2/2, \quad (\rho, \mathbf{v})(t, \mathbf{x}) \rightarrow (\rho(\tau, \mathbf{x}), \tau' \mathbf{v}(\tau, \mathbf{x})), \quad \tau' = \frac{d\tau}{dt} = t$$

$$\partial_t \rho + \operatorname{div}(\rho \mathbf{v}) = \mathbf{0}, \quad \partial_t(\rho \mathbf{v}) + \operatorname{div}(\rho \mathbf{v} \otimes \mathbf{v}) = -\operatorname{grad} p : \rightarrow$$

$$(\partial_\tau \rho + \operatorname{div}(\rho \mathbf{v}))\tau' = \mathbf{0}, \quad \tau'' \rho \mathbf{v} + (\tau')^2 [\partial_\tau(\rho \mathbf{v}) + \operatorname{div}(\rho \mathbf{v} \otimes \mathbf{v})] = -\operatorname{grad} p$$

By dropping the terms written in red* we indeed get the Darcy law

$$\rho \mathbf{v} = -\operatorname{grad} p, \quad \partial_\tau \rho = \Delta p$$

(*) motivation: $(\tau')^2 = 2\tau = t^2 \ll 1 = \tau''$ for short times.

A richer example: "Darcy's Magnetohydrodynamics"

$$\partial_t \rho + \operatorname{div}(\rho \mathbf{v}) = \mathbf{0}, \quad \rho \mathbf{v} = \operatorname{div}(\eta \mathbf{B} \otimes \mathbf{B}) - \operatorname{grad} \mathbf{p}$$

$$\partial_t \mathbf{B} + \operatorname{curl}(\mathbf{B} \times \mathbf{v}) + \operatorname{curl}(\mu \operatorname{curl}(\nu \mathbf{B})) = \mathbf{0}$$

Here $(\rho, \mathbf{p}, \mathbf{v}, \mathbf{B}) \in \mathbb{R}^{1+1+3+3}$ are the density, pressure, velocity and magnetic fields, $(\mu, \nu, \eta, \mathbf{p})$ being given functions of ρ .

A richer example: "Darcy's Magnetohydrodynamics"

$$\partial_t \rho + \operatorname{div}(\rho \mathbf{v}) = \mathbf{0}, \quad \rho \mathbf{v} = \operatorname{div}(\eta \mathbf{B} \otimes \mathbf{B}) - \operatorname{grad} \mathbf{p}$$

$$\partial_t \mathbf{B} + \operatorname{curl}(\mathbf{B} \times \mathbf{v}) + \operatorname{curl}(\mu \operatorname{curl}(\nu \mathbf{B})) = \mathbf{0}$$

Here $(\rho, \mathbf{p}, \mathbf{v}, \mathbf{B}) \in \mathbb{R}^{1+1+3+3}$ are the density, pressure, velocity and magnetic fields, $(\mu, \nu, \eta, \mathbf{p})$ being given functions of ρ .

As $\mathbf{B} = \mathbf{0}$, we recognize the classical Darcy law.

A richer example: "Darcy's Magnetohydrodynamics"

$$\partial_t \rho + \operatorname{div}(\rho \mathbf{v}) = \mathbf{0}, \quad \rho \mathbf{v} = \operatorname{div}(\eta \mathbf{B} \otimes \mathbf{B}) - \operatorname{grad} \mathbf{p}$$

$$\partial_t \mathbf{B} + \operatorname{curl}(\mathbf{B} \times \mathbf{v}) + \operatorname{curl}(\mu \operatorname{curl}(\nu \mathbf{B})) = \mathbf{0}$$

Here $(\rho, \mathbf{p}, \mathbf{v}, \mathbf{B}) \in \mathbb{R}^{1+1+3+3}$ are the density, pressure, velocity and magnetic fields, $(\mu, \nu, \eta, \mathbf{p})$ being given functions of ρ .

As $\mathbf{B} = \mathbf{0}$, we recognize the classical Darcy law.

A priori, such a model, combining Darcy law and MHD, is very far from "first principles" and must require many steps to be derived!

A direct derivation from... "pure" Geometry and Physics

Assume $\mu = \nu = \eta = \rho^{-1}$ **with pressure law** $\mathbf{p} = -\rho^{-1}$
("Chaplygin" pressure, used in Cosmology, with sound speed $(\frac{dp}{d\rho})^{1/2} = \rho^{-1}$)

A direct derivation from... "pure" Geometry and Physics

Assume $\mu = \nu = \eta = \rho^{-1}$ **with pressure law** $\mathbf{p} = -\rho^{-1}$
("Chaplygin" pressure, used in Cosmology, with sound speed $(\frac{dp}{d\rho})^{1/2} = \rho^{-1}$)

$$\partial_t \rho + \operatorname{div}(\rho \mathbf{v}) = \mathbf{0}, \quad \partial_t \mathbf{B} + \operatorname{curl}(\mathbf{B} \times \mathbf{v}) = -\operatorname{curl}(\rho^{-1} \operatorname{curl}(\rho^{-1} \mathbf{B}))$$

$$\rho \mathbf{v} = \operatorname{div}(\rho^{-1} \mathbf{B} \otimes \mathbf{B}) + \operatorname{grad}(\rho^{-1})$$

A direct derivation from... "pure" Geometry and Physics

Assume $\mu = \nu = \eta = \rho^{-1}$ with pressure law $\mathbf{p} = -\rho^{-1}$
("Chaplygin" pressure, used in Cosmology, with sound speed $(\frac{dp}{d\rho})^{1/2} = \rho^{-1}$)

$$\partial_t \rho + \operatorname{div}(\rho \mathbf{v}) = \mathbf{0}, \quad \partial_t \mathbf{B} + \operatorname{curl}(\mathbf{B} \times \mathbf{v}) = -\operatorname{curl}(\rho^{-1} \operatorname{curl}(\rho^{-1} \mathbf{B}))$$

$$\rho \mathbf{v} = \operatorname{div}(\rho^{-1} \mathbf{B} \otimes \mathbf{B}) + \operatorname{grad}(\rho^{-1})$$

This can be derived, through the simple "Darcy time-rescaling"
 $t \rightarrow t^2/2$, from a very "pure" equation from Geometry and Physics,

A direct derivation from... "pure" Geometry and Physics

Assume $\mu = \nu = \eta = \rho^{-1}$ **with pressure law** $\mathbf{p} = -\rho^{-1}$
("Chaplygin" pressure, used in Cosmology, with sound speed $(\frac{dp}{d\rho})^{1/2} = \rho^{-1}$)

$$\partial_t \rho + \operatorname{div}(\rho \mathbf{v}) = \mathbf{0}, \quad \partial_t \mathbf{B} + \operatorname{curl}(\mathbf{B} \times \mathbf{v}) = -\operatorname{curl}(\rho^{-1} \operatorname{curl}(\rho^{-1} \mathbf{B}))$$

$$\rho \mathbf{v} = \operatorname{div}(\rho^{-1} \mathbf{B} \otimes \mathbf{B}) + \operatorname{grad}(\rho^{-1})$$

This can be derived, through the simple "Darcy time-rescaling" $t \rightarrow t^2/2$, from a very "pure" equation from Geometry and Physics, designed by Born and Infeld in 1934 and used in... String Theory!



Max BORN (1882-1970) Physics Nobel prize 1954

Geometric definition of the Born-Infeld theory

In a $d + 1$ dimensional Lorentzian space-time manifold of metric $g_{ij}dx^i dx^j$ the BI theory involves closed 2-forms $\mathcal{B} = \mathcal{B}_{ij}dx^i \wedge dx^j$

Geometric definition of the Born-Infeld theory

In a $d + 1$ dimensional Lorentzian space-time manifold of metric $g_{ij}dx^i dx^j$ the BI theory involves closed 2-forms $\mathcal{B} = \mathcal{B}_{ij}dx^i \wedge dx^j$ that are critical points * of the fully covariant** action

$$\int (\sqrt{-\det \mathbf{g}} - \sqrt{-\det(\mathbf{g} + \mathcal{B})})$$

* for compactly supported variations/** invariant under changes of coordinates/

Geometric definition of the Born-Infeld theory

In a $d + 1$ dimensional Lorentzian space-time manifold of metric $g_{ij}dx^i dx^j$ the BI theory involves closed 2-forms $\mathcal{B} = \mathcal{B}_{ij}dx^i \wedge dx^j$ that are critical points * of the fully covariant** action

$$\int (\sqrt{-\det \mathbf{g}} - \sqrt{-\det(\mathbf{g} + \mathcal{B})})$$

* for compactly supported variations/** invariant under changes of coordinates/

We will concentrate on the 3+1 Minkowski space of special relativity

Geometric definition of the Born-Infeld theory

In a $d + 1$ dimensional Lorentzian space-time manifold of metric $g_{ij}dx^i dx^j$ the BI theory involves closed 2-forms $\mathcal{B} = \mathcal{B}_{ij}dx^i \wedge dx^j$ that are critical points * of the fully covariant** action

$$\int (\sqrt{-\det \mathbf{g}} - \sqrt{-\det(\mathbf{g} + \mathcal{B})})$$

* for compactly supported variations/** invariant under changes of coordinates/

We will concentrate on the 3+1 Minkowski space of special relativity (as Max Born and Leopold Infeld did in 1934).

Born-Infeld in traditional electromagnetic notations

After simple but tedious calculations, the Born-Infeld equations read, using classical electromagnetic notations,

$$\partial_t \mathbf{B} + \operatorname{curl} \left(\frac{\mathbf{B} \times (\mathbf{D} \times \mathbf{B}) + \mathbf{D}}{\sqrt{1 + \mathbf{D}^2 + \mathbf{B}^2 + (\mathbf{D} \times \mathbf{B})^2}} \right) = \mathbf{0}, \quad \operatorname{div} \mathbf{B} = 0$$

$$\partial_t \mathbf{D} + \operatorname{curl} \left(\frac{\mathbf{D} \times (\mathbf{D} \times \mathbf{B}) - \mathbf{B}}{\sqrt{1 + \mathbf{D}^2 + \mathbf{B}^2 + (\mathbf{D} \times \mathbf{B})^2}} \right) = \mathbf{0}, \quad \operatorname{div} \mathbf{D} = 0$$

Born-Infeld in traditional electromagnetic notations

After simple but tedious calculations, the Born-Infeld equations read, using classical electromagnetic notations,

$$\partial_t \mathbf{B} + \text{curl} \left(\frac{\mathbf{B} \times (\mathbf{D} \times \mathbf{B}) + \mathbf{D}}{\sqrt{1 + \mathbf{D}^2 + \mathbf{B}^2 + (\mathbf{D} \times \mathbf{B})^2}} \right) = \mathbf{0}, \quad \text{div} \mathbf{B} = 0$$

$$\partial_t \mathbf{D} + \text{curl} \left(\frac{\mathbf{D} \times (\mathbf{D} \times \mathbf{B}) - \mathbf{B}}{\sqrt{1 + \mathbf{D}^2 + \mathbf{B}^2 + (\mathbf{D} \times \mathbf{B})^2}} \right) = \mathbf{0}, \quad \text{div} \mathbf{D} = 0$$

They return the usual homogeneous Maxwell equations in the vacuum (terms in black) as \mathbf{B} , \mathbf{D} are fields or very small amplitude.

The "energy-momentum" conservation laws

By Noether's theorem*, we get 4 extra conservation laws

$$\partial_t \rho + \operatorname{div}(\rho \mathbf{v}) = \mathbf{0}, \quad \partial_t(\rho \mathbf{v}) + \operatorname{div}(\rho \mathbf{v} \otimes \mathbf{v} - \frac{\mathbf{B} \otimes \mathbf{B} - \mathbf{D} \otimes \mathbf{D}}{\rho}) = \operatorname{grad}(\rho^{-1})$$

$$\mathbf{v} = \frac{\mathbf{D} \times \mathbf{B}}{\rho}, \quad \rho = \sqrt{1 + \mathbf{D}^2 + \mathbf{B}^2 + (\mathbf{D} \times \mathbf{B})^2}$$

* which can be applied here since, once the Minkowski space-time is chosen, the Born-Infeld action gets invariant under time and space translations/

The augmented 10x10 Born-Infeld system

(following Y.B. ARMA 2004 and Y.B./Weinan Yong JMP 2006)

It is consistent (and much simpler) to consider $(\mathbf{B}, \mathbf{D}, \rho, \mathbf{v})$ as independent variables solving the 10×10 augmented system (which includes the 4 extra conservation laws)

$$\partial_t \mathbf{B} + \text{curl}(\mathbf{B} \times \mathbf{v} + \rho^{-1} \mathbf{D}) = \mathbf{0}, \quad \partial_t \mathbf{D} + \text{curl}(\mathbf{D} \times \mathbf{v} - \rho^{-1} \mathbf{B}) = \mathbf{0}$$

$$\partial_t \rho + \text{div}(\rho \mathbf{v}) = \mathbf{0}, \quad \partial_t(\rho \mathbf{v}) + \text{div}(\rho \mathbf{v} \otimes \mathbf{v} - \frac{\mathbf{B} \otimes \mathbf{B} - \mathbf{D} \otimes \mathbf{D}}{\rho}) = \text{grad}(\rho^{-1})$$

The augmented 10x10 Born-Infeld system

(following Y.B. ARMA 2004 and Y.B./Weinan Yong JMP 2006)

It is consistent (and much simpler) to consider $(\mathbf{B}, \mathbf{D}, \rho, \mathbf{v})$ as independent variables solving the 10×10 augmented system (which includes the 4 extra conservation laws)

$$\partial_t \mathbf{B} + \text{curl}(\mathbf{B} \times \mathbf{v} + \rho^{-1} \mathbf{D}) = \mathbf{0}, \quad \partial_t \mathbf{D} + \text{curl}(\mathbf{D} \times \mathbf{v} - \rho^{-1} \mathbf{B}) = \mathbf{0}$$

$$\partial_t \rho + \text{div}(\rho \mathbf{v}) = \mathbf{0}, \quad \partial_t(\rho \mathbf{v}) + \text{div}(\rho \mathbf{v} \otimes \mathbf{v} - \frac{\mathbf{B} \otimes \mathbf{B} - \mathbf{D} \otimes \mathbf{D}}{\rho}) = \text{grad}(\rho^{-1})$$

while ignoring the algebraic constraint $\mathbf{v} = \frac{\mathbf{D} \times \mathbf{B}}{\rho}$, $\rho = (1 + D^2 + B^2 + (D \times B)^2)^{1/2}$.

Darcy time-rescaling of the (augmented) BI equations

$$t \rightarrow \tau = t^2/2, \quad (\rho, B, D, v)(t, x) \rightarrow (\rho(\tau, x), B(\tau, x), \tau' D(\tau, x), \tau' v(\tau, x))$$

Darcy time-rescaling of the (augmented) BI equations

$$t \rightarrow \tau = t^2/2, \quad (\rho, \mathbf{B}, \mathbf{D}, \mathbf{v})(t, \mathbf{x}) \rightarrow (\rho(\tau, \mathbf{x}), \mathbf{B}(\tau, \mathbf{x}), \tau' \mathbf{D}(\tau, \mathbf{x}), \tau' \mathbf{v}(\tau, \mathbf{x}))$$

$$\rightarrow: \quad (\partial_\tau \rho + \operatorname{div}(\rho \mathbf{v}))\tau' = \mathbf{0}, \quad (\partial_\tau \mathbf{B} + \operatorname{curl}(\mathbf{B} \times \mathbf{v} + \rho^{-1} \mathbf{D}))\tau' = \mathbf{0}$$

Darcy time-rescaling of the (augmented) BI equations

$$t \rightarrow \tau = t^2/2, \quad (\rho, \mathbf{B}, \mathbf{D}, \mathbf{v})(t, \mathbf{x}) \rightarrow (\rho(\tau, \mathbf{x}), \mathbf{B}(\tau, \mathbf{x}), \tau' \mathbf{D}(\tau, \mathbf{x}), \tau' \mathbf{v}(\tau, \mathbf{x}))$$

$$\rightarrow: \quad (\partial_\tau \rho + \operatorname{div}(\rho \mathbf{v})) \tau' = \mathbf{0}, \quad (\partial_\tau \mathbf{B} + \operatorname{curl}(\mathbf{B} \times \mathbf{v} + \rho^{-1} \mathbf{D})) \tau' = \mathbf{0}$$

$$\rightarrow: \quad \tau'' \mathbf{D} + (\tau')^2 [\partial_\tau \mathbf{D} + \operatorname{curl}(\mathbf{D} \times \mathbf{v})] = \operatorname{curl}(\rho^{-1} \mathbf{B})$$

Darcy time-rescaling of the (augmented) BI equations

$$t \rightarrow \tau = t^2/2, \quad (\rho, \mathbf{B}, \mathbf{D}, \mathbf{v})(t, \mathbf{x}) \rightarrow (\rho(\tau, \mathbf{x}), \mathbf{B}(\tau, \mathbf{x}), \tau' \mathbf{D}(\tau, \mathbf{x}), \tau' \mathbf{v}(\tau, \mathbf{x}))$$

$$\rightarrow: \quad (\partial_\tau \rho + \operatorname{div}(\rho \mathbf{v})) \tau' = \mathbf{0}, \quad (\partial_\tau \mathbf{B} + \operatorname{curl}(\mathbf{B} \times \mathbf{v} + \rho^{-1} \mathbf{D})) \tau' = \mathbf{0}$$

$$\rightarrow: \quad \tau'' \mathbf{D} + (\tau')^2 [\partial_\tau \mathbf{D} + \operatorname{curl}(\mathbf{D} \times \mathbf{v})] = \operatorname{curl}(\rho^{-1} \mathbf{B})$$

$$\tau'' \rho \mathbf{v} + (\tau')^2 [\partial_\tau(\rho \mathbf{v}) + \operatorname{div}(\rho \mathbf{v} \otimes \mathbf{v} - \frac{\mathbf{D} \otimes \mathbf{D}}{\rho})] = \operatorname{div}(\frac{\mathbf{B} \otimes \mathbf{B}}{\rho}) + \operatorname{grad}(\rho^{-1})$$

Darcy time-rescaling of the (augmented) BI equations

$$t \rightarrow \tau = t^2/2, \quad (\rho, \mathbf{B}, \mathbf{D}, \mathbf{v})(t, \mathbf{x}) \rightarrow (\rho(\tau, \mathbf{x}), \mathbf{B}(\tau, \mathbf{x}), \tau' \mathbf{D}(\tau, \mathbf{x}), \tau' \mathbf{v}(\tau, \mathbf{x}))$$

$$\rightarrow: \quad (\partial_\tau \rho + \operatorname{div}(\rho \mathbf{v}))\tau' = \mathbf{0}, \quad (\partial_\tau \mathbf{B} + \operatorname{curl}(\mathbf{B} \times \mathbf{v} + \rho^{-1} \mathbf{D}))\tau' = \mathbf{0}$$

$$\rightarrow: \quad \tau'' \mathbf{D} + (\tau')^2 [\partial_\tau \mathbf{D} + \operatorname{curl}(\mathbf{D} \times \mathbf{v})] = \operatorname{curl}(\rho^{-1} \mathbf{B})$$

$$\tau'' \rho \mathbf{v} + (\tau')^2 [\partial_\tau (\rho \mathbf{v}) + \operatorname{div}(\rho \mathbf{v} \otimes \mathbf{v} - \frac{\mathbf{D} \otimes \mathbf{D}}{\rho})] = \operatorname{div}(\frac{\mathbf{B} \otimes \mathbf{B}}{\rho}) + \operatorname{grad}(\rho^{-1})$$

Ignoring the terms written in red leads to the desired model of Darcy MHD, with "constitutive laws" $\mu = \nu = \eta = -\rho = 1/\rho$!

Energy conservation versus energy dissipation

In the (augmented) BI equations, the total energy is *conserved*

$$\frac{d}{dt} \int \frac{1 + \mathbf{B}^2 + \mathbf{D}^2 + (\rho \mathbf{v})^2}{\rho} dx = 0$$

Energy conservation versus energy dissipation

In the (augmented) BI equations, the total energy is *conserved*

$$\frac{d}{dt} \int \frac{1 + \mathbf{B}^2 + \mathbf{D}^2 + (\rho\mathbf{v})^2}{\rho} dx = 0$$

For the rescaled version, we find

$$\tau' \frac{d}{d\tau} \int \frac{1 + \mathbf{B}^2}{\rho} dx + \tau'^3 \frac{d}{d\tau} \int \frac{\mathbf{D}^2 + (\rho\mathbf{v})^2}{\rho} dx = -2\tau'' \tau' \int \frac{\mathbf{D}^2 + (\rho\mathbf{v})^2}{\rho} dx$$

Energy conservation versus energy dissipation

In the (augmented) BI equations, the total energy is *conserved*

$$\frac{d}{dt} \int \frac{1 + \mathbf{B}^2 + \mathbf{D}^2 + (\rho\mathbf{v})^2}{\rho} dx = 0$$

For the rescaled version, we find

$$\tau' \frac{d}{d\tau} \int \frac{1 + \mathbf{B}^2}{\rho} dx + \tau'^3 \frac{d}{d\tau} \int \frac{\mathbf{D}^2 + (\rho\mathbf{v})^2}{\rho} dx = -2\tau'' \tau' \int \frac{\mathbf{D}^2 + (\rho\mathbf{v})^2}{\rho} dx$$

After dropping the red terms, we get the energy *dissipation* relation

$$\frac{d}{d\tau} \int \frac{1 + \mathbf{B}^2}{\rho} dx = -2 \int \frac{\mathbf{D}^2 + (\rho\mathbf{v})^2}{\rho} dx$$

(remember: $(\tau')^2 = 2\tau = t^2 \ll 1 = \tau''$)

So, we have seen "Darcy Magnetohydrodynamics" as an example of a theory, apparently very far remote from very "first principles",

So, we have seen "Darcy Magnetohydrodynamics" as an example of a theory, apparently very far remote from very "first principles", which, nevertheless, at least for specific constitutive laws, can be derived, at once and directly, from a very "pure" geometric theory.

So, we have seen "Darcy Magnetohydrodynamics" as an example of a theory, apparently very far remote from very "first principles", which, nevertheless, at least for specific constitutive laws, can be derived, at once and directly, from a very "pure" geometric theory.

Thus, the gap between "pure" and "applied" theories may be much narrower than usually admitted!

So, we have seen "Darcy Magnetohydrodynamics" as an example of a theory, apparently very far remote from very "first principles", which, nevertheless, at least for specific constitutive laws, can be derived, at once and directly, from a very "pure" geometric theory.

Thus, the gap between "pure" and "applied" theories may be much narrower than usually admitted!

MERCI DE VOTRE ATTENTION ET SURTOUT...

So, we have seen "Darcy Magnetohydrodynamics" as an example of a theory, apparently very far remote from very "first principles", which, nevertheless, at least for specific constitutive laws, can be derived, at once and directly, from a very "pure" geometric theory.

Thus, the gap between "pure" and "applied" theories may be much narrower than usually admitted!

**MERCI DE VOTRE ATTENTION ET SURTOUT...
MERCI ET FELICITATIONS A JEAN ET JEROME!**